

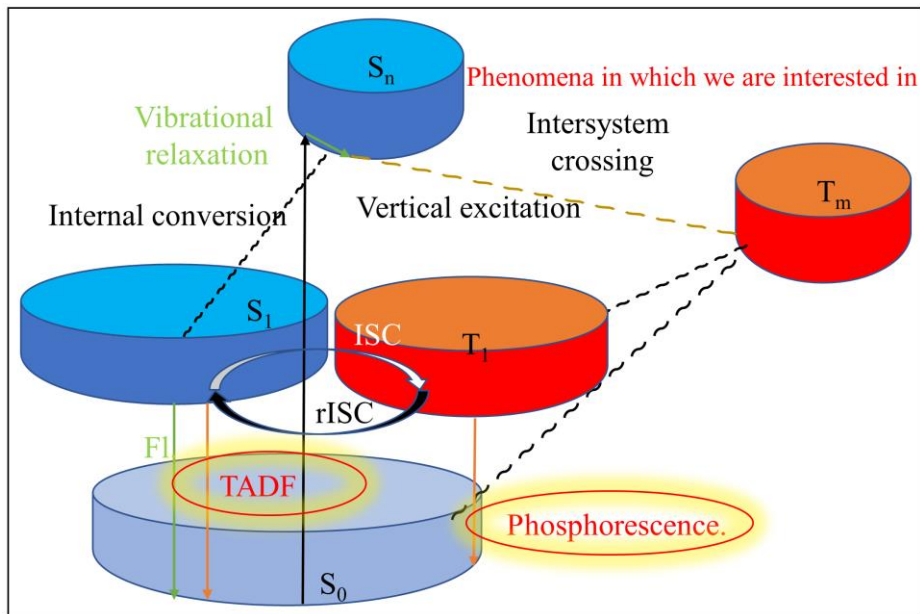
Intersystem Crossing in Boron-Based Donor–Spiro–Acceptor Organic Chromophore: A Detailed Theoretical Study

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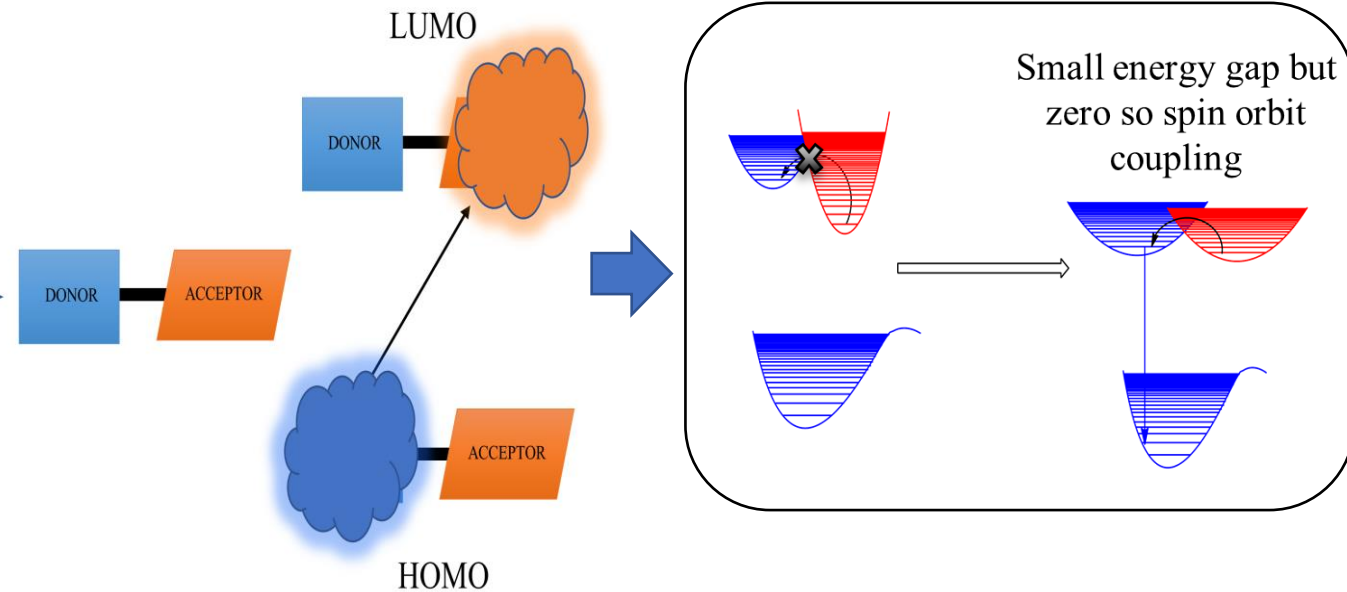
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INTRODUCTION:

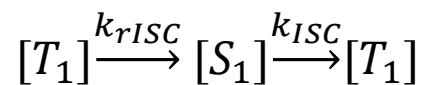
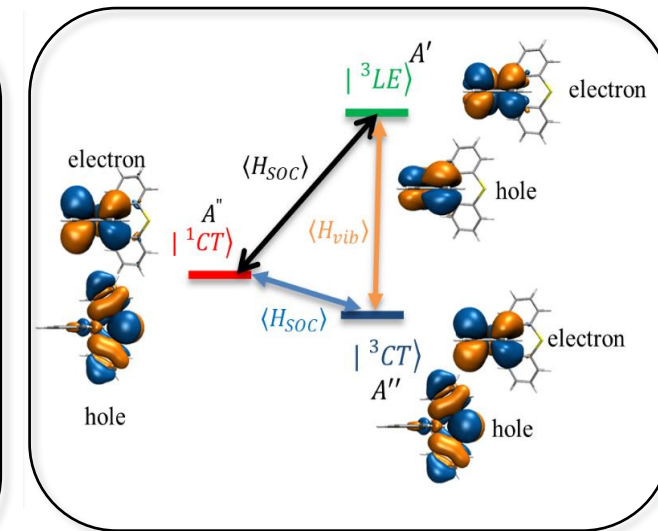
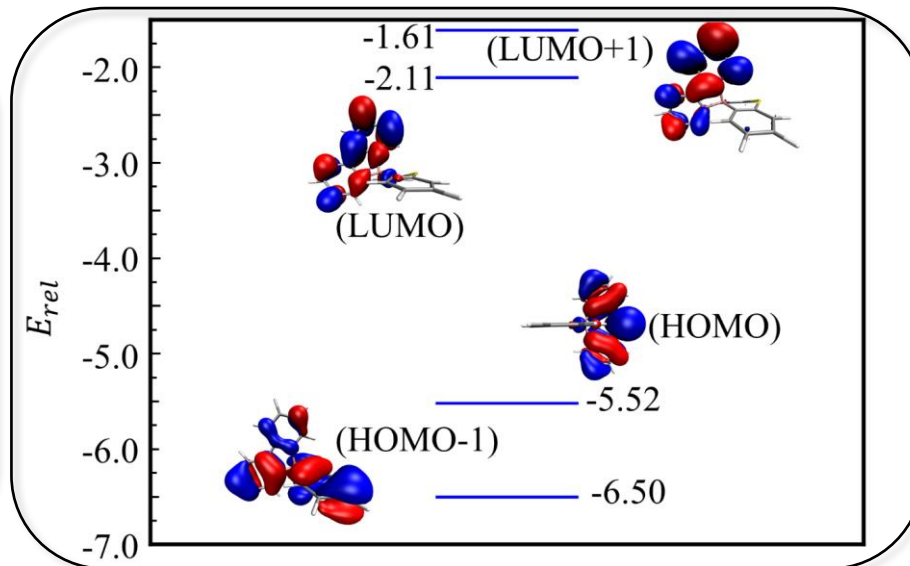
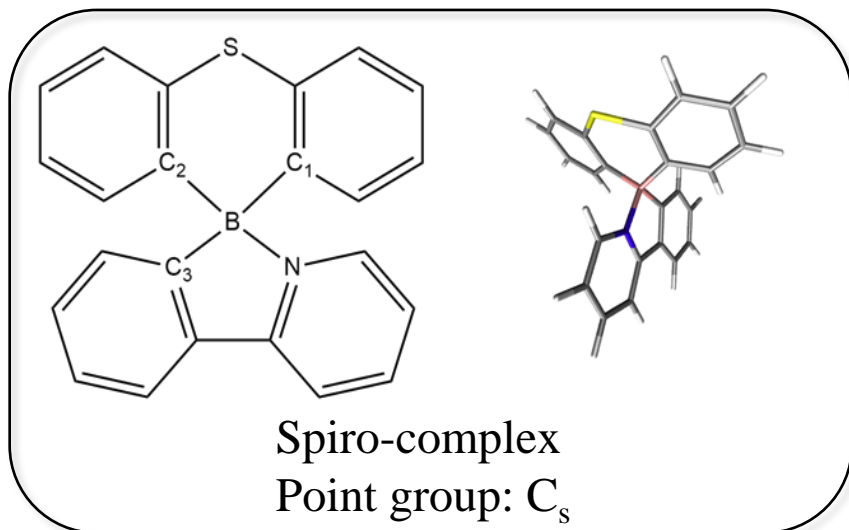


Design Strategy



- Stanoppi, M.; Lorbach, A. Boron-based donor-spiro-acceptor compounds exhibiting thermally activated delayed fluorescence (TADF). Dalton Transactions 2018, 47, 10394–10398.
- Gibson, J.; Monkman, A. P.; Penfold, T. J. The importance of vibronic coupling for efficient reverse intersystem crossing in thermally activated delayed fluorescence molecules. ChemPhysChem 2016, 17, 2956.

RESULTS AND DISCUSSIONS



So, S_1 state population at time t is given by the following expression,

$$[S_1] = [T_1]_0 - \frac{[T_1]_0}{(k_{rISC} + k_{ISC})} (k_{ISC} + k_{rISC} \times e^{-(k_{ISC} + k_{rISC})t})$$

Reorganization energy	k_{ISC}	k_{rISC}
(λ_M)		
$\lambda_M = 0.3 \text{ eV}$	$2.45 \times 10^5 \text{ s}^{-1}$	$9.70 \times 10^4 \text{ s}^{-1}$
$\lambda_M = 0.4 \text{ eV}$	$8.12 \times 10^4 \text{ s}^{-1}$	$3.21 \times 10^4 \text{ s}^{-1}$

$[T_1]_0 = 100.0$ as the initial population, we get, $[S_1] = 0.0003$ after 100 ps

Non-radiative decay from n th Singlet state to m th Triplet state due to ISC can be calculated using Fermi-Golden rule

$$k_{ISC}^{nm} = \frac{2\pi}{\hbar} \rho_{FC} |\langle S_n | \widehat{H}_{SO} | T_m \rangle|^2$$

RESULTS AND DISCUSSIONS

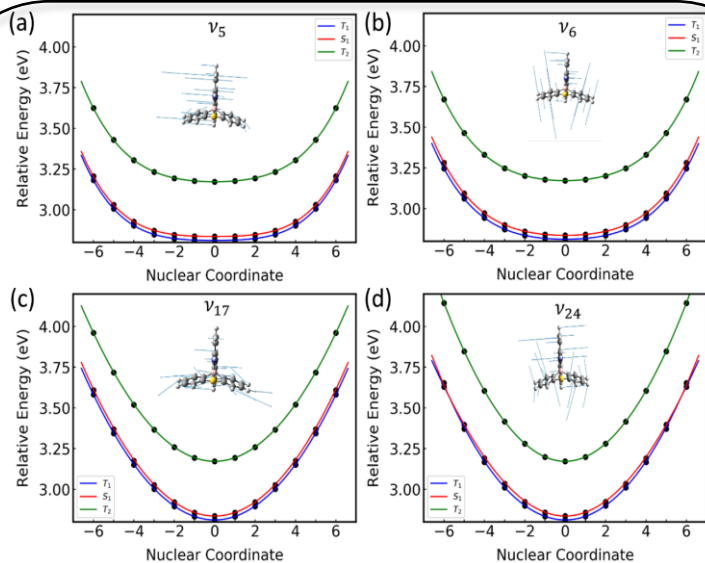
Quantum dynamics simulation: Spin-Vibronic Coupling Hamiltonian

$$\begin{aligned}
 & \left(\begin{array}{l} E_{T_1}(Q_0) + \sum_i \frac{\omega_i}{2} \partial_{Q_i}^2 + \frac{\omega_i}{2} Q_i^2 + \frac{\gamma^{T_1}}{2} Q_i^2 + \frac{\varepsilon^{T_1}}{24} Q_i^4 \\ E_{SOC}^{S1T1} \\ \sum_i \lambda_i Q_i \end{array} \right) \quad \begin{array}{l} E_{SOC}^{S1T1} \\ E_{SOC}^{S1T2} \\ \sum_i \lambda_i Q_i \end{array} \\
 & \left(\begin{array}{l} E_{S_1}(Q_0) + \sum_i \frac{\omega_i}{2} \partial_{Q_i}^2 + \frac{\omega_i}{2} Q_i^2 + \frac{\gamma^{S_1}}{2} Q_i^2 + \frac{\varepsilon^{S_1}}{24} Q_i^4 \\ E_{SOC}^{S1T2} \\ \sum_i \lambda_i Q_i \end{array} \right) \quad \begin{array}{l} E_{SOC}^{S1T2} \\ E_{T_2}(Q_0) + \sum_i \frac{\omega_i}{2} \partial_{Q_i}^2 + \frac{\omega_i}{2} Q_i^2 + \frac{\gamma^{T_2}}{2} Q_i^2 + \frac{\varepsilon^{T_2}}{24} Q_i^4 \end{array}
 \end{aligned}$$

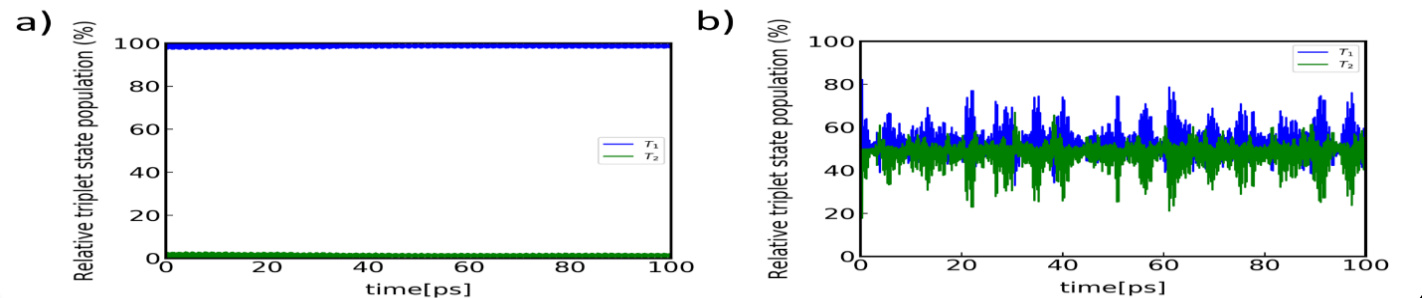
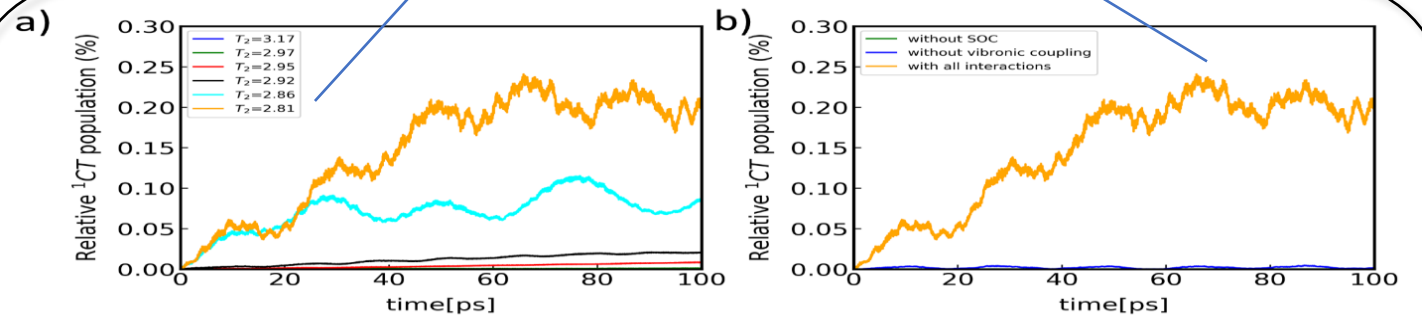
$$\begin{aligned}
 H_{nn} &= E_n(Q_0) + \sum_i \frac{\omega_i}{2} \partial_{Q_i}^2 + \frac{\omega_i}{2} Q_i^2 + \frac{\gamma^n}{2} Q_i^2 + \frac{\varepsilon^n}{24} Q_i^4 \\
 H_{nm} &= \sum_i \lambda_i^{nm} Q_i
 \end{aligned}$$

$$\lambda_i^{nm} = \sqrt{\frac{1}{8} \frac{\partial^2}{\partial Q_i^2} [V_m(Q) - V_n(Q)]^2}$$

- We modelled this Hamiltonian to do excited state dynamics
- Inter state vibronic coupling (λ) parameters are calculated from potential energy scan
- Only low energy vibrations are taken into account
 - When both the triplet states have same energy
 - At that energy level both SOCME and vibronic coupling are operative



Low frequency vibrations



CONCLUSIONS

- Spiro arrangement of D-A system leads to CT nature in both S_1 and T_1 transitions
- SOCME between these two states is 0.078 cm^{-1}
- T_2 is LE in nature and it helps in rISC through spin-vibronic coupling (2^{nd} order perturbation)
- Apart from ΔE_{ST} , ΔE_{TT} gap is also playing a bigger role in rISC

COMPUTATIONAL METHODOLOGY

- ❖ Excited State energy potential scan were performed using B3LYP functional with 6-31+g(d,p) basis set using Gaussian 16 package.
- ❖ SCOME calculations were done using ZORA Hamiltonian with ADF2017 package using B3LYP functional and TZP basis set.
- ❖ Excited State dynamics were performed using MCTDH.

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THANK YOU